Control and Cybernetics

vol. **42** (2013) No. 2

A generalized fractional calculus of variations^{*}

by

Tatiana Odzijewicz¹, Agnieszka B. Malinowska² and Delfim F. M. $Torres^1$

¹CIDMA — Center for Research and Development in Mathematics and Applications, Department of Mathematics, University of Aveiro, 3810-193 Aveiro, Portugal tatianao@ua.pt; delfim@ua.pt ²Faculty of Computer Science, Bialystok University of Technology, 15-351 Białystok, Poland a.malinowska@pb.edu.pl

Abstract: We study incommensurate fractional variational problems in terms of a generalized fractional integral with Lagrangians depending on classical derivatives and generalized fractional integrals and derivatives. We obtain necessary optimality conditions for the basic and isoperimetric problems, transversality conditions for free boundary value problems, and a generalized Noether type theorem.

Keywords: generalized fractional operators; fractional variational analysis; Euler–Lagrange equations; natural boundary conditions; Noether's theorem; damped harmonic oscillator.

1. Introduction

Till recently, it was believed that Lagrangian and Hamiltonian mechanics were not valid in the presence of nonconservative forces such as friction (Lanezos, 1970). In the last years, however, several approaches have been investigated in order to find a Lagrangian or a Hamiltonian description for classes of dissipative (or dissipative-looking) systems (Cieśliński and Nikiciuk, 2010; Crampin et al., 2010; Kobe et al., 1986; Musielak, 2008; Pradeep et al., 2009). One way to have a Lagrangian and a Hamiltonian formulation, for both conservative and nonconservative systems, was proposed by Fred Riewe in 1996 and consists in using fractional derivatives (Riewe, 1996, 1997). Riewe's papers (Riewe, 1996, 1997) gave rise to a new and important research field, called *the fractional calculus of variations* (Malinowska and Torres, 2012). Nowadays the subject is of

^{*}Submitted: October 2012; Accepted: April 2013.

strong interest, and many results of variational analysis were extended to the noninteger case (see, e.g., Almeida et al. 2010, 2012; Almeida and Torres, 2009; Bastos et al., 2011a,b; Cresson, 2007; Frederico and Torres, 2008; Mozyrska and Torres, 2011; Odzijewicz et al., 2012a; Odzijewicz and Torres, 2011). Here we study problems of calculus of variations with generalized fractional operators (Agrawal, 2010; Odzijewicz et al., 2012b,c). Generalized fractional integrals are given as a linear combination of left and right fractional integrals with general kernels. Generalized fractional Riemann-Liouville and Caputo derivatives are defined as a composition of classical derivatives and generalized fractional integrals. In the first problem, we ask how to determine the extremizers of a functional defined by a generalized fractional integral involving n generalized fractional Caputo derivatives and n generalized fractional integrals. All these operators have different (noninteger) orders. We obtain necessary optimality conditions, and in the case of free boundary values, also natural boundary conditions. Next, we derive Euler-Lagrange type equations for an extended isoperimetric problem and we obtain a Noether type theorem.

The text is organized as follows. In Section 2 we give the definitions and main properties of the generalized fractional operators. We prove Euler-Lagrange equations for the fundamental generalized problem in Section 3, and natural boundary conditions for free boundary value problems in Section 4. Section 5 is devoted to the generalized isoperimetric problem and Section 6 to Noether's theorem. Finally, in Section 7 we present an application of our results to the damped harmonic oscillator.

2. Preliminaries

We start by defining the generalized fractional operators (Agrawal, 2010). As particular cases, by choosing appropriate kernels, such operators are reduced to the standard fractional integrals and derivatives of fractional calculus (see, e.g., Kilbas et al., 2006; Klimek, 2009; Podlubny, 1999). Throughout the text, α denotes a real number between zero and one. Following Almeida et al. (2012), we use round brackets for the arguments of functions, and square brackets for the arguments of operators.

DEFINITION 1 (The generalized fractional integral). The operator K_P^{α} is given by

$$K_P^{\alpha}[f](x) := K_P^{\alpha}[t \mapsto f(t)](x) = p \int_a^x k_{\alpha}(x,t)f(t)dt + q \int_x^b k_{\alpha}(t,x)f(t)dt,$$

where $P = \langle a, x, b, p, q \rangle$ is the parameter set (*p*-set for brevity), $x \in [a, b]$, p, qare real numbers, and $k_{\alpha}(x, t)$ is a kernel which may depend on α . The operator K_P^{α} is referred as the operator K (K-op for simplicity) of order α and *p*-set P.

Note that if we define

$$G(x,t) := \begin{cases} pk_{\alpha}(x,t) & \text{if } t < x, \\ qk_{\alpha}(t,x) & \text{if } t \ge x, \end{cases}$$

then the operator K_P^{α} can be written in the form

$$K_P^{\alpha}[f](x) = K_P^{\alpha}[t \mapsto f(t)](x) = \int_a^b G(x,t)f(t)dt.$$

Thus, the generalized fractional integral is a Fredholm operator, one of the oldest and most respectable class of operators that arise in the theory of integral equations (Helemskii, 2006; Polyanin and Manzhirov, 1998).

EXAMPLE 1. 1. Let $k_{\alpha}(t-\tau) = \frac{1}{\Gamma(\alpha)}(t-\tau)^{\alpha-1}$ and $0 < \alpha < 1$. If $P = \langle a, t, b, 1, 0 \rangle$, then

$$K_P^{\alpha}[f](t) = \frac{1}{\Gamma(\alpha)} \int_a^t (t-\tau)^{\alpha-1} f(\tau) d\tau =: {}_a I_t^{\alpha}[f](t)$$

is the left Riemann–Liouville fractional integral of order $\alpha;$ if $P=\langle a,t,b,0,1\rangle,$ then

$$K_P^{\alpha}[f](t) = \frac{1}{\Gamma(\alpha)} \int_t^b (\tau - t)^{\alpha - 1} f(\tau) d\tau =: {}_t I_b^{\alpha}[f](t)$$

is the right Riemann-Liouville fractional integral of order α . 2. For $k_{\alpha}(t-\tau) = \frac{1}{\Gamma(\alpha(t,\tau))}(t-\tau)^{\alpha(t,\tau)-1}$ and $P = \langle a, t, b, 1, 0 \rangle$

$$K_P^{\alpha}[f](t) = \int_a^t \frac{1}{\Gamma(\alpha(t,\tau)} (t-\tau)^{\alpha(t,\tau)-1} f(\tau) d\tau =: {_aI_t^{\alpha(\cdot,\cdot)}[f](t)}$$

is the left Riemann-Liouville fractional integral of variable order $\alpha(t,\tau)$, and for $P = \langle a, t, b, 0, 1 \rangle$

$$K_P^{\alpha}[f](t) = \int_t^b \frac{1}{\Gamma(\alpha(\tau, t))} (\tau - t)^{\alpha(t, \tau) - 1} f(\tau) d\tau =: {}_t I_b^{\alpha(\cdot, \cdot)}[f](t)$$

is the right Riemann-Liouville fractional integral of variable order $\alpha(t,\tau)$ (Odzijewicz et al., 2013).

3. For $0 < \alpha < 1$, $k_{\alpha}(t,\tau) = \frac{1}{\Gamma(\alpha)} \left(\log \frac{t}{\tau} \right)^{\alpha-1} \frac{1}{\tau}$ and $P = \langle a, t, b, 1, 0 \rangle$, the operator K_P^{α} reduces to the left Hadamard fractional integral (Pooseh et al., 2012),

$$K_P^{\alpha}[f](t) = \frac{1}{\Gamma(\alpha)} \int_a^t \left(\log \frac{t}{\tau} \right)^{\alpha - 1} \frac{f(\tau) d\tau}{\tau} =: {_aJ_t^{\alpha}[f](t)},$$

and for $P = \langle a, t, b, 0, 1 \rangle$ operator K_P reduces to the right Hadamard fractional integral,

$$K_P^{\alpha}[f](t) = \frac{1}{\Gamma(\alpha)} \int_t^b \left(\log \frac{\tau}{t}\right)^{\alpha - 1} \frac{f(\tau)d\tau}{\tau} =: {}_t J_b^{\alpha}[f](t).$$

 Generalized fractional integrals can be also reduced to, e.g., Riesz, Katugampola or Kilbas fractional operators. Their definitions can be found in (Katugampola, 2010; Kilbas and Saigo, 2004; Kilbas et al., 2006).

Next results yield boundedness of the generalized fractional integral.

THEOREM 1 (see Example 6 of Helemskii, 2006). Let $\alpha \in (0,1)$ and $P = \langle a, x, b, p, q \rangle$. If k_{α} is a square integrable function on the square $\Delta = [a, b] \times [a, b]$, then $K_P^{\alpha} : L_2([a, b]) \to L_2([a, b])$ is well defined, linear, and bounded operator.

THEOREM 2 (see Odzijewicz et al., 2012b,c). Let $k_{\alpha} \in L_1([0, b-a])$ be a difference kernel, that is, $k_{\alpha}(x,t) = k_{\alpha}(x-t)$. Then, $K_P^{\alpha} : L_1([a,b]) \to L_1([a,b])$ is a well defined bounded and linear operator.

THEOREM 3 (see Theorem 2.4 of Odzijewicz et al., 2012c). Let $P = \langle a, x, b, p, q \rangle$. If $k_{1-\alpha}$ is a difference kernel, $k_{1-\alpha} \in L_1([0, b-a])$ and $f \in AC([a, b])$, then $K_P^{1-\alpha}[f]$ belongs to AC([a, b]).

The generalized fractional derivatives A_P^{α} and B_P^{α} are defined in terms of the generalized fractional integral K-op.

DEFINITION 2 (Generalized Riemann–Liouville fractional derivative). Let P be a given parameter set and $0 < \alpha < 1$. The operator A_P^{α} is defined by $A_P^{\alpha} := D \circ K_P^{1-\alpha}$, where D denotes the standard derivative operator, and is referred to as the operator A (A-op) of order α and p-set P.

REMARK 1. Operator A is well-defined for all functions f such that $K_P^{1-\alpha}[f]$ is differentiable. Theorem 3 ensures that the domain of A is nonempty.

DEFINITION 3 (Generalized Caputo fractional derivative). Let P be a given parameter set and $\alpha \in (0,1)$. The operator B_P^{α} is defined by $B_P^{\alpha} := K_P^{1-\alpha} \circ D$, where D denotes the standard derivative operator, and is referred to as the operator B (B-op) of order α and p-set P.

REMARK 2. Operator B is well-defined for differentiable functions.

EXAMPLE 2. The standard Riemann-Liouville and Caputo fractional derivatives (see, e.g., Kilbas et al., 2006; Podlubny, 1999; Klimek, 2009) are easily obtained from the general kernel operators A_P^{α} and B_P^{α} , respectively. Let $k_{\alpha}(t-\tau) = \frac{1}{\Gamma(1-\alpha)}(t-\tau)^{-\alpha}$, $\alpha \in (0,1)$. If $P = \langle a, t, b, 1, 0 \rangle$, then

$$A_P^{\alpha}[f](t) = \frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_a^t (t-\tau)^{-\alpha} f(\tau) d\tau =: {}_a D_t^{\alpha}[f](t)$$

is the standard left Riemann-Liouville fractional derivative of order α , while

$$B_P^{\alpha}[f](t) = \frac{1}{\Gamma(1-\alpha)} \int_{a}^{t} (t-\tau)^{-\alpha} f'(\tau) d\tau =: {}_{a}^{C} D_t^{\alpha}[f](t)$$

is the standard left Caputo fractional derivative of order $\alpha;$ if $P=\langle a,t,b,0,1\rangle,$ then

$$-A_P^{\alpha}[f](t) = -\frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_t^b (\tau-t)^{-\alpha} f(\tau) d\tau =: {}_t D_b^{\alpha}[f](t)$$

is the standard right Riemann-Liouville fractional derivative of order α , while

$$-B_{P}^{\alpha}[f](t) = -\frac{1}{\Gamma(1-\alpha)} \int_{t}^{b} (\tau-t)^{-\alpha} f'(\tau) d\tau =: {}_{t}^{C} D_{b}^{\alpha}[f](t)$$

is the standard right Caputo fractional derivative of order α .

The following theorems give integration by parts formulas for operators A, B and K. For detailed proofs we refer the reader to (Odzijewicz et al., 2012b,c).

THEOREM 4. Let $\alpha \in (0,1)$, $P = \langle a,t,b,p,q \rangle$, k_{α} be a square-integrable function on $\Delta = [a,b] \times [a,b]$, and $f,g \in L_2([a,b])$. The generalized fractional integral K_P^{α} satisfies the integration by parts formula

$$\int_{a}^{b} g(x) K_{P}^{\alpha}[f](x) dx = \int_{a}^{b} f(x) K_{P^{*}}^{\alpha}[g](x) dx, \qquad (1)$$

where $P^* = \langle a, t, b, q, p \rangle$.

THEOREM 5. Let $\alpha \in (0,1)$, $P = \langle a,t,b,p,q \rangle$, and k_{α} be a square integrable function on $\Delta = [a,b] \times [a,b]$. If functions $f, K_{P^*}^{1-\alpha}[g] \in AC([a,b])$, then

$$\int_{a}^{b} g(x) B_{P}^{\alpha}[f](x) dx = f(x) K_{P^{*}}^{1-\alpha}[g](x) \Big|_{a}^{b} - \int_{a}^{b} f(x) A_{P^{*}}^{\alpha}[g](x) dx, \qquad (2)$$

where $P^* = \langle a, t, b, q, p \rangle$.

THEOREM 6. Let $0 < \alpha < 1$, $P = \langle a, x, b, p, q \rangle$, and k_{α} be a difference kernel such that $k_{\alpha} \in L_1(0, b-a)$. If $f \in L_1([a, b])$ and $g \in C([a, b])$, then the operator K_P^{α} satisfies the integration by parts formula (1).

THEOREM 7. Let $\alpha \in (0,1)$, $P = \langle a,t,b,p,q \rangle$, and $k_{\alpha} \in L_1((0,b-a])$ be a difference kernel. If functions $f, g \in AC([a,b])$, then formula (2) holds.



For $\mathbf{f} = [f_1, \ldots, f_N] : [a, b] \to \mathbb{R}^N$, where $N \in \mathbb{N}$, we put

 $A_{P}^{\alpha}[\mathbf{f}](x) := [A_{P}^{\alpha}[f_{1}](x), \dots, A_{P}^{\alpha}[f_{N}](x)],$ $B_{P}^{\alpha}[\mathbf{f}](x) := [B_{P}^{\alpha}[f_{1}](x), \dots, B_{P}^{\alpha}[f_{N}](x)],$ $K_{P}^{\alpha}[\mathbf{f}](x) := [K_{P}^{\alpha}[f_{1}](x), \dots, K_{P}^{\alpha}[f_{N}](x)].$

3. The generalized fundamental variational problem

We consider the problem of finding a function $\mathbf{y} = [y_1, \ldots, y_N]$ that gives an extremum (minimum or maximum) to the functional

$$\mathcal{J}(\mathbf{y}) = K_P^{\alpha} \Big[t \mapsto F\big(t, \mathbf{y}(t), \mathbf{y}'(t), B_{P_1}^{\beta_1} [\mathbf{y}](t), \dots, B_{P_n}^{\beta_n} [\mathbf{y}](t), K_{R_1}^{\gamma_1} [\mathbf{y}](t), \dots, K_{R_m}^{\gamma_m} [\mathbf{y}](t) \big) \Big](b) \quad (3)$$

subject to the boundary conditions

$$\mathbf{y}(a) = \mathbf{y}_a, \quad \mathbf{y}(b) = \mathbf{y}_b, \tag{4}$$

where $\alpha, \beta_i, \gamma_k \in (0, 1), P = \langle a, b, b, 1, 0 \rangle, P_i = \langle a, t, b, p_i, q_i \rangle$, and $R_k = \langle a, t, b, r_k, s_k \rangle$, $i = 1, \ldots, n, k = 1, \ldots, m$. For simplicity of notation, we introduce the operator $\{\cdot\}_{P_D, R_I}^{\beta, \gamma}$ defined by

$$\{\mathbf{y}\}_{P_D,R_I}^{\beta,\gamma}(t) := \left(t, \mathbf{y}(t), \mathbf{y}'(t), B_{P_D}^{\beta}\left[\tau \mapsto \mathbf{y}(\tau)\right](t), K_{R_I}^{\gamma}\left[\tau \mapsto \mathbf{y}(\tau)\right](t)\right),$$

where

$$B_{P_D}^{\beta} := \left(B_{P_1}^{\beta_1}, \dots, B_{P_n}^{\beta_n} \right), \quad K_{R_I}^{\gamma} := \left(K_{R_1}^{\gamma_1}, \dots, K_{R_m}^{\gamma_m} \right).$$

The operator K_P^{α} has kernel $k_{\alpha}(x,t)$ and, for i = 1, ..., n and k = 1, ..., m, operators $B_{P_i}^{\beta_i}$ and $K_{R_k}^{\gamma_k}$ have kernels $h_{1-\beta_i}(t,\tau)$ and $h_{\gamma_k}(t,\tau)$, respectively. In the sequel we assume that:

(H1) the Lagrangian $F \in C^1([a,b] \times \mathbb{R}^{N \times (n+m+2)}; \mathbb{R});$

(H2) functions $t \mapsto k_{\alpha}(b,t)\partial_j F\{\mathbf{y}\}_{P_D,R_I}^{\beta,\gamma}(t)$ and

$$D\left[t \mapsto \partial_{N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t)k_{\alpha}(b,t)\right],$$

$$A_{P_{i}^{*}}^{\beta_{i}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(i+1)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right],$$

$$K_{R_{k}^{*}}^{\gamma_{k}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(n+1+k)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right]$$

are continuous on (a, b), j = 2, ..., N + 1, i = 1, ..., n, k = 1, ..., m; (H3) functions $t \mapsto \partial_{N+j} F \{\mathbf{y}\}_{P_D, R_I}^{\beta, \gamma}(t) k_{\alpha}(b, t)$ and

$$K_{P_{i}^{*}}^{1-\beta_{i}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(i+1)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right]$$

belong to AC([a, b]), j = 2, ..., N + 1, i = 1, ..., n;

(H4) for i = 1, ..., n, k = 1, ..., m, the kernels $k_{\alpha}(x, t), h_{1-\beta_i}(t, \tau)$ and $h_{\gamma_k}(t, \tau)$ are such that we are able to use Theorems 4, 5, 6 and/or 7.

DEFINITION 4. A function $\mathbf{y} \in C^1([a,b]; \mathbb{R}^N)$ is said to be admissible for the fractional variational problem (3)-(4) if functions $B_{P_i}^{\beta_i}[\mathbf{y}]$ and $K_{R_k}^{\gamma_k}[\mathbf{y}]$, $i = 1, \ldots, n, k = 1, \ldots, m$ exist and are continuous on the interval [a,b], and \mathbf{y} satisfies the boundary conditions (4).

THEOREM 8. If \mathbf{y} is a solution to problem (3)–(4), then \mathbf{y} satisfies the system of generalized Euler-Lagrange equations

$$k_{\alpha}(b,t)\partial_{j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}\left(t\right) - \sum_{i=1}^{n} A_{P_{i}^{*}}^{\beta_{i}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(i+1)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}\left(\tau\right)\right]\left(t\right) + \sum_{k=1}^{m} K_{R_{k}^{*}}^{\gamma_{k}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(n+1+k)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}\left(\tau\right)\right]\left(t\right)$$

$$- \frac{d}{dt}\left(\partial_{N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}\left(t\right)k_{\alpha}(b,t)\right) = 0$$

$$(5)$$

for all $t \in (a, b), j = 2, ..., N + 1$.

Proof. The proof is analogous to that in Odzijewicz et al. (2012b, Theorem 4.2).

4. Generalized free-boundary variational problem

Assume now that in problem (3)–(4) the boundary conditions (4) are substituted by

$$\mathbf{y}(a)$$
 is free and $\mathbf{y}(b) = \mathbf{y}_b$. (6)

THEOREM 9. If \mathbf{y} is a solution to the problem of extremizing functional (3) with (6) as the boundary conditions, then \mathbf{y} satisfies the system of Euler-Lagrange equations (5). Moreover, the extra system of natural boundary conditions

$$\partial_{N+j} F\left\{\mathbf{y}\right\}_{P_{D,R_{I}}}^{\beta,\gamma}(a) k_{\alpha}(b,a) + \sum_{i=1}^{n} K_{P_{i}^{*}}^{1-\beta_{i}}\left[\tau \mapsto \partial_{(i+1)N+j} F\left\{\mathbf{y}\right\}_{P_{D,R_{I}}}^{\beta,\gamma}(\tau) k_{\alpha}(b,\tau)\right](a) = 0, \quad (7)$$

j = 2, ..., N + 1, holds.

Proof. The proof is analogous to that of Odzijewicz et al. (2012b, Theorem 5.1).

5. Generalized isoperimetric problem

Let $\xi \in \mathbb{R}$. Among all functions $\mathbf{y} : [a, b] \to \mathbb{R}^N$ satisfying the boundary conditions

$$\mathbf{y}(a) = \mathbf{y}_a, \quad \mathbf{y}(b) = \mathbf{y}_b, \tag{8}$$

and an isoperimetric constraint of the form

$$\mathcal{I}(\mathbf{y}) = K_P^{\alpha} \left[G\left\{ \mathbf{y} \right\}_{P_D, R_I}^{\beta, \gamma} \right](b) = \xi, \tag{9}$$

we look for those that extremize (i.e., minimize or maximize) the functional

$$\mathcal{J}(\mathbf{y}) = K_P^{\alpha} \left[F \left\{ \mathbf{y} \right\}_{P_D, R_I}^{\beta, \gamma} \right] (b).$$
(10)

For i = 1, ..., n, k = 1, ..., m, operators K_P^{α} , $B_{P_i}^{\beta_i}$ and $K_{R_k}^{\gamma_k}$, as well as function F, are the same as in problem (3)–(4). Moreover, we assume that functional (9) satisfies hypotheses (H1)–(H4).

DEFINITION 5. A function $\mathbf{y} : [a, b] \to \mathbb{R}^N$ is said to be admissible for problem (8)–(10) if functions $B_{P_i}^{\beta_i}[\mathbf{y}]$ and $K_{R_k}^{\gamma_k}[\mathbf{y}]$, $i = 1, \ldots, n$, $k = 1, \ldots, m$, exist and are continuous on [a, b], and \mathbf{y} satisfies the boundary conditions (8) and the isoperimetric constraint (9).

DEFINITION 6. An admissible function $\mathbf{y} \in C^1([a, b], \mathbb{R}^N)$ is said to be an extremal for \mathcal{I} if it satisfies the system of Euler-Lagrange equations (5) associated with functional in (9).

THEOREM 10. If \mathbf{y} is a solution to the isoperimetric problem (8)–(10) and is not an extremal for \mathcal{I} , then there exists a real constant λ such that

$$k_{\alpha}(b,t)\partial_{j}H\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t) + \sum_{k=1}^{m} K_{R_{k}^{k}}^{\gamma_{k}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(n+1+k)N+j}H\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right](t)$$
$$-\sum_{i=1}^{n} A_{P_{i}^{i}}^{\beta_{i}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(i+1)N+j}H\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right](t)$$
$$-\frac{d}{dt}\left(\partial_{j+N}H\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t)k_{\alpha}(b,t)\right) = 0$$

for all $t \in (a, b)$, j = 2, ..., N+1, where $P_i^* = \langle a, t, b, q_i, p_i \rangle$, $R_k^* = \langle a, t, b, s_k, r_k \rangle$, and

$$H(t, y, u, v, w) = F(t, y, u, v, w) - \lambda G(t, y, u, v, w).$$

Proof. The proof is analogous to that of Odzijewicz et al. (2012b, Theorem 6.3).

6. Generalized fractional Noether's theorem

Emmy Noether's theorem on extremal functionals, establishing that certain symmetries imply conservation laws (constants of motion), has been called "the most important theorem in physics since the Pythagorean theorem". For a recent account of Noether's theorem and possible applications in physics, from many different points of view, we refer the reader to Neuenschwander (2011). Formulations in the more general context of optimal control can be found in Gouveia et al. (2006), Torres (2002). Conservation laws appear naturally in closed systems. In presence of nonconservative or dissipative forces, the constants of motion are broken and Noether's classical theorem ceases to be valid. It is still possible, however, to obtain Noether type theorems that cover both conservative and nonconservative cases. Roughly speaking, one can prove that Noether's conservation laws are still valid if a new term, involving the nonconservative forces, is added to the standard constants of motion (Frederico and Torres, 2007a). The first Noether theorem for the fractional calculus of variations was obtained in 2007 (Frederico and Torres, 2007b). Since then, the subject attracted a lot of attention. The state of the art is given in the book Malinowska and Torres (2012). Here we obtain a Noether theorem for generalized fractional variational problems.

DEFINITION 7. We say that the functional (3) is invariant under an ε -parameter group of infinitesimal transformations

$$\hat{\mathbf{y}}(t) = \mathbf{y}(t) + \varepsilon \boldsymbol{\xi}(t, \mathbf{y}(t)) + o(\varepsilon)$$
(11)

if for any subinterval $[t_a, t_b] \subseteq [a, b]$ one has

$$K_{\bar{P}}^{\alpha}\left[t\mapsto F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t)\right](t_{b})=K_{\bar{P}}^{\alpha}\left[t\mapsto F\left\{\hat{\mathbf{y}}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t)\right](t_{b}),\tag{12}$$

where $\bar{P} = \langle t_a, t_b, t_b, 1, 0 \rangle$.

THEOREM 11. If functional (3) is invariant under an ε -parameter group of infinitesimal transformations, then

$$\sum_{j=2}^{N+1} \left(\partial_j F \left\{ \mathbf{y} \right\}_{P_D,R_I}^{\beta,\gamma}(t) \cdot \xi_{j-1}(t,\mathbf{y}(t)) + \partial_{N+j} F \left\{ \mathbf{y} \right\}_{P_D,R_I}^{\beta,\gamma}(t) \cdot \frac{d}{dt} \xi_{j-1}(t,\mathbf{y}(t)) + \sum_{i=1}^n \partial_{(i+1)N+j} F \left\{ \mathbf{y} \right\}_{P_D,R_I}^{\beta,\gamma}(t) \cdot B_{P_i}^{\beta_i}[\tau \mapsto \xi_{j-1}(\tau,\mathbf{y}(\tau))](t) + \sum_{k=1}^m \partial_{(n+1+k)N+j} F \left\{ \mathbf{y} \right\}_{P_D,R_I}^{\beta,\gamma}(t) \cdot K_{R_i}^{\gamma_i}[\tau \mapsto \xi_{j-1}(\tau,\mathbf{y}(\tau))](t) = 0.$$
(13)

Proof. Since, by hypothesis, condition (12) is satisfied for any subinterval $[t_a, t_b] \subseteq [a, b]$, we have

$$F\{\mathbf{y}\}_{P_{D},R_{I}}^{\beta,\gamma}(t) = F\{\hat{\mathbf{y}}\}_{P_{D},R_{I}}^{\beta,\gamma}(t).$$
(14)



Differentiating (14) with respect to ε , then putting $\varepsilon = 0$, and applying definitions and properties of generalized fractional operators, we obtain (13).

In order to state the Noether theorem in a compact form, we introduce the following operators:

$$\mathbf{D}_{P}^{\alpha}[f,g](t) := \frac{1}{k_{\alpha}(b,t)} f(t) \cdot A_{P^{*}}^{\alpha}[g](t) + g(t) \cdot B_{P}^{\alpha}[f](t),$$
(15)

$$\mathbf{I}_{P}^{\alpha}[f,g](t) := \frac{-1}{k_{\alpha}(b,t)} f(t) \cdot K_{P^{*}}^{\alpha}[g](t) + g(t) \cdot K_{P}^{\alpha}[f](t),$$
(16)

where P^* denotes the dual *p*-set of *P*, that is, if $P = \langle a, t, b, p, q \rangle$, then $P^* = \langle a, t, b, q, p \rangle$.

THEOREM 12 (Generalized fractional Noether theorem). If functional (3) is invariant under an ε -parameter group of infinitesimal transformations (11), then

$$\sum_{j=2}^{N+1} \left(\sum_{i=1}^{n} \mathbf{D}_{P_{i}}^{\beta_{i}} \left[\tau \mapsto \xi_{j-1}(\tau, \mathbf{y}(\tau)), \tau \mapsto k_{\alpha}(b, \tau) \partial_{(i+1)N+j} F\left\{\mathbf{y}\right\}_{P_{D}, R_{I}}^{\beta, \gamma}(\tau) \right] (t) \\ + \sum_{k=1}^{m} \mathbf{I}_{R_{k}}^{\gamma_{k}} \left[\tau \mapsto \xi_{j-1}(\tau, \mathbf{y}(\tau)), \tau \mapsto k_{\alpha}(b, \tau) \partial_{(n+1+k)N+j} F\left\{\mathbf{y}\right\}_{P_{D}, R_{I}}^{\beta, \gamma}(\tau) \right] (t) \\ + \frac{d}{dt} \left(\xi_{j-1}\left(t, \mathbf{y}(t)\right) \cdot \partial_{N+j} F\left\{\mathbf{y}\right\}_{P_{D}, R_{I}}^{\beta, \gamma}(t) \right) \\ + \xi_{j-1}(t, \mathbf{y}(t)) \cdot \partial_{N+j} F\left\{\mathbf{y}\right\}_{P_{D}, R_{I}}^{\beta, \gamma}(t) \cdot \frac{1}{k_{\alpha}(b, t)} \frac{d}{dt} k_{\alpha}(b, t) \right) = 0$$

$$(17)$$

for any generalized fractional extremal \mathbf{y} of \mathcal{J} and for all $t \in (a, b)$.

Proof. By Theorem 8 we have

$$k_{\alpha}(b,t)\partial_{j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t) = \sum_{i=1}^{n} A_{P_{i}^{*}}^{\beta_{i}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(i+1)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right](t)$$
$$-\sum_{k=1}^{m} K_{R_{i}^{*}}^{\gamma_{k}}\left[\tau \mapsto k_{\alpha}(b,\tau)\partial_{(n+1+k)N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau)\right](t)$$
$$+\frac{d}{dt}\left(\partial_{N+j}F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t)k_{\alpha}(b,t)\right)$$
(18)

for all $t \in (a, b)$, j = 2, ..., N + 1. Substituting (18) into (13), we obtain

$$\begin{split} \sum_{j=2}^{N+1} \left[\frac{1}{k_{\alpha}(b,t)} \cdot \xi_{j-1}(t,\mathbf{y}(t)) \left(\sum_{i=1}^{n} A_{P_{i}^{*}}^{\beta_{i}} \left[\tau \mapsto k_{\alpha}(b,\tau) \partial_{(i+1)N+j} F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau) \right](t) \right. \\ \left. \left. - \sum_{k=1}^{m} K_{R_{i}^{*}}^{\gamma_{k}} \left[\tau \mapsto k_{\alpha}(b,\tau) \partial_{(n+1+k)N+j} F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(\tau) \right](t) \right. \\ \left. \left. + \frac{d}{dt} \left(\partial_{N+j} F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t) k_{\alpha}(b,t) \right) \right) + \partial_{N+j} F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t) \cdot \frac{d}{dt} \xi_{j-1}\left(t,\mathbf{y}(t)\right) \right. \\ \left. \left. + \sum_{i=1}^{n} \partial_{(i+1)N+j} F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t) \cdot B_{P_{i}}^{\beta_{i}}[\tau \mapsto \xi_{j-1}\left(\tau,\mathbf{y}(\tau)\right)](t) \right. \\ \left. \left. + \sum_{k=1}^{m} \partial_{(n+1+k)N+j} F\left\{\mathbf{y}\right\}_{P_{D},R_{I}}^{\beta,\gamma}(t) \cdot K_{R_{i}}^{\gamma_{i}}[\tau \mapsto \xi_{j-1}(\tau,\mathbf{y}(\tau))](t) \right] = 0. \end{split}$$

Finally, we arrive at (17) by (15) and (16).

EXAMPLE 3. Let $P = \langle a, t, b, p, q \rangle$. Consider the following problem:

$$\mathcal{J}[y] = \int_{a}^{b} F(t, B_{P}^{\alpha}[y](t)) dt \longrightarrow \min$$

$$y(a) = y_{a}, \quad y(b) = y_{b},$$
(19)

 $and \ transformations$

$$\hat{y}(t) = y(t) + \varepsilon c + o(\varepsilon),$$
(20)

where c is a constant. For any $[t_a, t_b] \subseteq [a, b]$ we have

$$\int_{t_a}^{t_b} F(t, B_P^{\alpha}[y](t)) \, dt = \int_{t_a}^{t_b} F(t, B_P^{\alpha}[\hat{y}](t)) \, dt.$$

Therefore, $\mathcal{J}[y]$ is invariant under (20) and Theorem 12 asserts that

$$A_{P*}^{\alpha}[\tau \to \partial_2 F(\tau, B_P^{\alpha}[y](\tau))](t) = 0$$
(21)

along any generalized fractional extremal y. Note that equation (21) can be written in the form

$$\frac{d}{dt}\left(K_{P^*}^{\alpha}[\tau \to \partial_2 F\left(\tau, B_P^{\alpha}[y](\tau)\right)](t)\right) = 0.$$

In analogy with the classical approach, quantity $K_{P^*}^{\alpha}[\tau \to \partial_2 F(\tau, B_P^{\alpha}[y](\tau))](t)$ is called a generalized fractional constant of motion.

7. Applications to physics

If the functional (3) does not depend on B-ops and K-ops, then Theorem 8 gives the following result: if **y** is a solution to the problem of extremizing

$$\mathcal{J}(\mathbf{y}) = \int_{a}^{b} L\left(t, \mathbf{y}(t), \mathbf{y}'(t)\right) k_{\alpha}(b, t) dt$$
(22)

subject to $\mathbf{y}(a) = \mathbf{y}_a$ and $\mathbf{y}(b) = \mathbf{y}_b$, where $\alpha \in (0, 1)$, then

$$\partial_{j}L(t, \mathbf{y}(t), \mathbf{y}'(t)) - \frac{d}{dt}\partial_{N+j}L(t, \mathbf{y}(t), \mathbf{y}'(t)) = \frac{1}{k_{\alpha}(b, t)} \cdot \frac{d}{dt}k_{\alpha}(b, t)\partial_{N+j}L(t, \mathbf{y}(t), \mathbf{y}'(t)), \quad (23)$$

 $j = 2, \ldots, N + 1$. In addition, if we assume that functional (22) is invariant under transformations (11), then Noether's theorem yields that

$$\sum_{j=2}^{N+1} \left(\frac{d}{dt} \left(\xi_{j-1}(t, \mathbf{y}(t)) \cdot \partial_{N+j} L(t, \mathbf{y}(t), \mathbf{y}'(t)) \right) + \xi_{j-1}(t, \mathbf{y}(t)) \cdot \partial_{N+j} L(t, \mathbf{y}(t), \mathbf{y}'(t)) \cdot \frac{1}{k_{\alpha}(b, t)} \frac{d}{dt} k_{\alpha}(b, t) \right) = 0,$$

along any extremal of (22). Let us consider kernel $k_{\alpha}(b,t) = e^{\alpha(b-t)}$ and the Lagrangian for a three dimensional system:

$$L(\mathbf{y}, \dot{\mathbf{y}}) = \frac{1}{2}m(\dot{y_1}^2 + \dot{y_2}^2 + \dot{y_3}^2) - V(\mathbf{y}),$$

where $V(\mathbf{y})$ is the potential energy and m stands for the mass. Observe that an explicitly time dependent integrand $\tilde{L} = e^{\alpha(b-t)}L$ of functional (22) is known in the literature as the Bateman–Caldirola–Kanai (BCK) Lagrangian of a quantum dissipative system (Menon et al., 1997; Ghosh et al., 2009). But in our case the Lagrangian of the system is L and not $e^{\alpha(b-t)}L$. The Euler–Lagrange equations (23) give the following system of second order ordinary differential equations:

$$\begin{cases} \ddot{y}_{1}(t) - \alpha \dot{y}_{1}(t) = -\frac{1}{m} \partial_{1} V(\mathbf{y}(t)) \\ \ddot{y}_{2}(t) - \alpha \dot{y}_{2}(t) = -\frac{1}{m} \partial_{2} V(\mathbf{y}(t)) \\ \ddot{y}_{3}(t) - \alpha \dot{y}_{3}(t) = -\frac{1}{m} \partial_{3} V(\mathbf{y}(t)). \end{cases}$$

If $\gamma := -\alpha$, then

$$\ddot{y}_i + \gamma \dot{y}_i + \frac{1}{m} \frac{\partial V}{\partial y_i} = 0, \tag{24}$$

i = 1, 2, 3, which are equations for the damped motion of a three-dimensional particle under the action of a force $\left[-\frac{\partial V}{\partial y_1}, -\frac{\partial V}{\partial y_2}, -\frac{\partial V}{\partial y_3}\right]$ (see, e.g., Herrera et al.,

1986). Choosing $V := k \frac{y_1^2 + y_2^2 + y_3^2}{2}$, we can transform (24) into equations for a damped simple harmonic oscillator:

$$\ddot{y}_i(t) + \gamma \dot{y}_i(t) + \omega^2 y_i(t) = 0,$$

i = 1, 2, 3, with $\omega^2 = \frac{k}{m}$. Now, let us consider the following Lagrangian:

$$L(\mathbf{y}, \dot{\mathbf{y}}) = \frac{1}{2}m\left(\dot{y_1}^2 + \dot{y_2}^2 + \dot{y_3}^2\right) - mgy_3^2.$$
(25)

We see at once that the Lagrangian (25) is invariant under the transformation

 $\hat{y}_1 = y_1 + \varepsilon, \ \hat{y}_2 = y_2, \ \hat{y}_3 = y_3.$

In this case Noether's theorem gives

$$\frac{d}{dt}(m\dot{y}_1) = \alpha m\dot{y}_1. \tag{26}$$

If $\alpha = 0$, then there is no friction and (26) yields the classical conservation of linear momentum $p_1 = m\dot{y}_1 = const$. Observe that the generalized momentum conjugate to y_i is $p_i = \frac{\partial L}{\partial \dot{y}_i} = m\dot{y}_i$, i = 1, 2, 3. This is not the case for the BCK Lagrangian (Ghosh et al., 2009), where the canonical momentum for y_i is $\tilde{p}_i = e^{\alpha(b-t)}m\dot{y}_i$, i = 1, 2, 3, that is different from the kinetic momentum. Now, let us suppose that L is variationally invariant under the transformation

$$\hat{y}_1 = y_1 \cos \varepsilon + y_2 \sin \varepsilon, \quad \hat{y}_2 = -y_1 \sin \varepsilon + y_2 \cos \varepsilon, \quad \hat{y}_3 = y_3.$$

Then $\xi_1 = y_2$, $\xi_2 = -y_1$ and $\xi_3 = 0$. For this case Noether's theorem yields

$$\frac{d}{dt}(m\dot{y}_1y_2 - my_1\dot{y}_2) - \alpha m(\dot{y}_1y_2 - y_1\dot{y}_2) = 0.$$
(27)

Note that for $\alpha = 0$ relation (27) gives the standard conservation law $p_1y_2 - p_2y_1 = const$ yielded by the classical Noether theorem (van Brunt, 2004, Section 9.3).

Acknowledgements

This work was supported by *FEDER* funds through *COMPETE* — Operational Programme Factors of Competitiveness ("Programa Operacional Factores de Competitividade") and by Portuguese funds through the *Center for Research and Development in Mathematics and Applications* (University of Aveiro) and the Portuguese Foundation for Science and Technology ("FCT — Fundação para a Ciência e a Tecnologia"), within project PEst-C/MAT/UI4106/2011 with COMPETE number FCOMP-01-0124-FEDER-022690. Tatiana Odzijewicz was also supported by FCT through the Ph.D. fellowship SFRH/BD/33865/2009; Agnieszka B. Malinowska by Bialystok University of Technology grant S/WI/02/2011; and Delfim F.M. Torres by FCT through the project PTDC/MAT/113470/2009.

The authors are very grateful to two anonymous referees for their valuable comments and helpful suggestions.

References

- AGRAWAL, O. P. (2010) Generalized variational problems and Euler-Lagrange equations. Comput. Math. Appl. 59 (5), 1852–1864.
- ALMEIDA, R., MALINOWSKA, A. B. and TORRES, D. F. M. (2010) A fractional calculus of variations for multiple integrals with application to vibrating string. J. Math. Phys. 51 (3), 033503.
- ALMEIDA, R., MALINOWSKA, A. B. and TORRES, D. F. M. (2012) Fractional Euler-Lagrange differential equations via Caputo derivatives. *Fractional Dynamics and Control*, Springer New York, Part 2, 109–118.
- ALMEIDA R. and TORRES D. F. M. (2009) Calculus of variations with fractional derivatives and fractional integrals. Appl. Math. Lett. 22 (12), 1816– 1820.
- BASTOS N. R. O., FERREIRA R. A. C. and TORRES, D. F. M. (2011a) Necessary optimality conditions for fractional difference problems of the calculus of variations. *Discrete Contin. Dyn. Syst.* 29 (2), 417–437.
- BASTOS N. R. O., FERREIRA R. A. C. and TORRES, D. F. M. (2011b) Discrete-time fractional variational problems. *Signal Process.* 91 (3), 513– 524.
- CIEŚLIŃSKI, J. L. AND NIKICIUK, T. (2010) A direct approach to the construction of standard and non-standard Lagrangians for dissipative-like dynamical systems with variable coefficients. J. Phys. A 43 (17), 175205.
- CRAMPIN, M., MESTDAG, T. and SARLET, W. (2010) On the generalized Helmholtz conditions for Lagrangian systems with dissipative forces. ZAMM Z. Angew. Math. Mech. 90 (6), 502–508.
- CRESSON, J. (2007) Fractional embedding of differential operators and Lagrangian systems. J. Math. Phys. 48 (3), 033504.
- FREDERICO, G. S. F. and TORRES, D. F. M. (2007a) Nonconservative Noether's theorem in optimal control. Int. J. Tomogr. Stat. 5 (W07), 109– 114.
- FREDERICO, G. S. F. and TORRES, D. F. M. (2007b) A formulation of Noether's theorem for fractional problems of the calculus of variations. J. Math. Anal. Appl. 334 (2), 834–846.
- FREDERICO, G. S. F. and TORRES, D. F. M. (2008) Fractional optimal control in the sense of Caputo and the fractional Noether's theorem. Int. Math. Forum 3 (10), 479–493.
- GHOSH, S., CHOUDHURI, A. and TALUKDAR, B. (2009) On the quantization of damped harmonic oscillator. Acta Phys. Polon. B 40 (1), 49–57.
- GIAQUINTA, M. and HILDEBRANDT, S. (1996) Calculus of Variations. I. Springer, Berlin.
- GOUVEIA, P. D. F., TORRES, D. F. M. and ROCHA, E. A. M. (2006) Symbolic computation of variational symmetries in optimal control. *Control Cybernet.* **35** (4), 831–849.
- HELEMSKII, A. YA. (2006) Lectures and Excercises on Functional Analysis. American Mathematical Society.

	Α	generalized	fractional	calculus	\mathbf{of}	variations
--	---	-------------	------------	----------	---------------	------------

- HERRERA, L., NÚÑEZ, L., PATIÑO, A. and RAGO, H. (1986) A variational principle and the classical and quantum mechanics of the damped harmonic oscillator. Am. J. Phys. 54 (3), 273–277.
- KATUGAMPOLA, U. N. (2010) New approach to a generalized fractional integral. Appl. Math. Comput. 218 (3), 860–865.
- KILBAS, A. A. and SAIGO, M. (2004) Generalized Mittag-Leffler function and generalized fractional calculus operators. *Integral Transform. Spec. Func.* 15 (1), 31–49.
- KILBAS, A. A., SRUVASTAVA, H. M. and TRUJILLO, J. J. (2006) Theory and Applications of Fractional Differential Equations. Elsevier, Amsterdam.
- KLIMEK, M. (2009) On solutions of linear fractional differential equations of a variational type. The Publishing Office of Czestochowa University of Technology, Czestochowa.
- KOBE, D. H., REALI, G. and SIENIUTYCZ, S. (1986) Lagrangians for dissipative systems. Am. J. Phys. 54, 997–999.
- LANEZOS, C. (1970) The Variational Principles of Mechanics. 4th edition, Dover, New York.
- MALINOWSKA, A. B. and TORRES, D. F. M. (2012) Introduction to the Fractional Calculus of Variations. Imp. Coll. Press, London.
- MENON, V. J., CHANANA, N. and SINGH, Y. (1997) A Fresh Look at the BCK Frictional Lagrangian. *Prog. Theor. Phys.* **98** (2), 321–329.
- MOZYRSKA, D. and TORRES, D. F. M. (2011) Modified optimal energy and initial memory of fractional continuous-time linear systems. *Signal Process.* **91** (3), 379–385.
- MUSIELAK, Z. E. (2008) Standard and non-standard Lagrangians for dissipative dynamical systems with variable coefficients. J. Phys. A 41 (5), 055205
- NEUENSCHWANDER, D.E. (2011) Emmy Noether's Wonderful Theorem. Johns Hopkins University Press, Baltimore.
- ODZIJEWICZ, T., MALINOWSKA, A.B. and TORRES, D. F. M. (2012a) Fractional variational calculus with classical and combined Caputo derivatives. *Nonlinear Anal.* **75** (3), 1507–1515.
- ODZIJEWICZ, T., MALINOWSKA, A.B. and TORRES, D. F. M. (2012b) Fractional calculus of variations in terms of a generalized fractional integral with applications to Physics. *Abstr. Appl. Anal.* 2012, Art. ID 871912.
- ODZIJEWICZ, T., MALINOWSKA, A.B. and TORRES, D. F. M. (2012c) Generalized fractional calculus with applications to the calculus of variations. *Comput. Math. Appl.* 64 (10), 3351–3366.
- ODZIJEWICZ, T., MALINOWSKA, A.B. and TORRES, D. F. M. (2013) Fractional variational calculus of variable order. In: A. Almeida, L. Castro, F.-O. Speck, Advances in Harmonic Analysis and Operator Theory, The Stefan Samko Anniversary Volume. Operator Theory: Advances and Applications 229, 291–301.
- ODZIJEWICZ, T. and TORRES, D.F.M. (2011) Fractional calculus of variations for double integrals. Balkan J. Geom. Appl. 16 (2), 102–113.

- PODLUBNY, I. (1999) Fractional Differential Equations. Academic Press, San Diego, CA.
- POLYANIN, A.D. and MANZHIROV, A.V. (1998) Handbook of Integral Equations. CRC, Boca Raton, FL.
- POOSEH, S., ALMEIDA, R. and TORRES, D. F. M. (2012) Expansion formulas in terms of integer-order derivatives for the Hadamard fractional integral and derivative. *Numer. Funct. Anal. Optim.* 33 (3), 301–319.
- PRADEEP, R. G., CHANDRASEKAR, V. K., SENTHILVELAN, M. and LAKSHMANAN, M. (2009) Nonstandard conserved Hamiltonian structures in dissipative/damped systems: nonlinear generalizations of damped harmonic oscillator. J. Math. Phys. 50 (5), 052901.
- RIEWE, F. (1996) Nonconservative Lagrangian and Hamiltonian mechanics. *Phys. Rev. E* **53** (2), 1890–1899.
- RIEWE, F. (1997) Mechanics with fractional derivatives. *Phys. Rev. E* 55 (3), part B, 3581–3592.
- TORRES, D. F. M. (2002) On the Noether theorem for optimal control. *Eur. J. Control* 8 (1), 56–63.
- VAN BRUNT, B. (2004) The Calculus of Variations. Universitext, Springer, New York.